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TITLE: On the relative energy of a static centrally symmetric gravitational field

PERIODICAL: Moscow. Universitet. Vestnik, Seriya. III. Fizika, astronomiya, no. 6, 1962, 45 - 55

TEXT: Energy and momentum of a static centrally symmetric gravitational field with respect to the center of symmetry of the system are calculated by a method devised by the author (Yu. A. Rylov. Vestn. Mosk. un-ta, ser. fiziki, astronomii, no. 5, 1962) from the relativistic field of gravitation $Q_{\beta\gamma}^{\alpha} = \delta_{\beta\gamma}^{\alpha}(x) - \Gamma_{\beta\gamma}^{\alpha}(x, x')$ (1). $\delta_{\beta\gamma}^{\alpha}$ are the Christoffel symbols in space-time V_4 in a certain system of coordinates K , $\Gamma_{\beta\gamma}^{\alpha}(x, x')$ are the Christoffel symbols in plane space E_x . These two spaces are tangent at the point x' of the coordinate system K_x . This relativistic field of gravitation is a two-point tensor describing the gravitational field at Card 1/4

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the point x with reference to the field at the point x' . First the world function corresponding to the live element $ds^2 = e^{\lambda} dt^2 - r^2(d\theta^2 + \sin^2\theta dy^2) - e^{\lambda} dr^2$ (11) is calculated, where $e^{\lambda} = 1 - 2\xi$, $e^{\lambda} = (1 - 2\xi)^{-1}$, $\xi = \alpha/2r$. α is the radius of gravitation, matter is assumed to be in the region $r < R$. Then the relativistic gravitational field

$$Q_{00}^0 = -\frac{2\xi}{r} \left(d - p + q - \frac{3}{2} \right), \quad Q_{13}^1 = \frac{\xi}{r} (d + 2p - 1),$$

$$Q_{22}^2 = \frac{\xi}{r} (d + 2p - 1), \quad Q_{00}^2 = -\frac{N-1}{r} \left(1 + N\xi - \frac{N}{N-1}\xi \right), \quad (28)$$

$$Q_{11}^3 = \xi(d + 2p + 1)r \sin^2\theta, \quad Q_{22}^3 = r\xi(d + 2p + 1).$$

is calculated by means of the transfer tensor

$$P_{\alpha}^{\beta} = -g_{\alpha\sigma}(x') G^{\sigma\beta}, \quad P_{\beta}^{\alpha} = -g^{\alpha\sigma}(x') G_{\sigma\beta},$$

(9). Finally, the four-momentum

$$P_{\beta} = \lim_{r \rightarrow \infty} \int_0^{2\pi} d\varphi \int_0^{\pi} \Lambda H_{\beta}^{\alpha} r^2 \sin\theta d\theta.$$

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(29) is obtained by means of

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$$P_p = \int H_p^{01} \sqrt{-D_x} d\sigma_{01} = \int \Delta H_p^{01} \sqrt{-g} d\sigma_{01}, \quad (5) \text{ and}$$

$$\Delta H_p^{01} = \frac{1}{2\pi} (\delta_p^0 (Q_{:,1}^0 - Q_{:,1}^{'0}) - \delta_p^1 (Q_{:,0}^0 - Q_{:,0}^{'0}) + Q_{:,0}^1 - Q_{:,1}^0), \quad (6), \text{ where}$$

$$H_0^{03} = \frac{1}{h'} H_0^{03},$$

$$H_1^{03} = -\frac{\sin \varphi}{r \sin \theta} H_1^{03} + \frac{\cos \theta \cos \varphi}{r} H_2^{03} + \sin \theta \cos \varphi H_3^{03},$$

$$H_2^{03} = \frac{\cos \varphi}{r \sin \theta} H_1^{03} + \frac{\cos \theta \sin \varphi}{r} H_2^{03} + \sin \theta \sin \varphi H_3^{03},$$

$$H_3^{03} = -\frac{\sin \theta}{r} H_2^{03} + \cos \theta H_3^{03}.$$

(30). (5) is integrated over a sphere of radius r at constant t ; therefore it is only necessary to know the asymptotic behavior of $Q_{\beta\gamma}^{\alpha}(\vec{r}, t, \vec{r}', t')$ at $t=t'=0$, $r'=0$, $r \rightarrow \infty$. It is shown that the energy of the gravitational field in Card 3/4

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relation to the point $x' = 0$ is negative.

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